# INTERPRETATION OF CIRRUS CLOUD PROPERTIES USING COINCIDENT SATELLITE AND LIDAR DATA DURING THE FIRE CIRRUS IFO

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#### 1. Introduction

The First ISCCP Regional Experiment (FIRE) Cirrus Intensive Field Observations (IFO) provide a unique opportunity to examine the relationships between the satellite-observed radiances and various parameters which describe the bulk properties of clouds, such as cloud amount and cloud-top height. In this paper, lidar-derived cloud altitude data, radiosonde data, and satellite-observed radiances are used to examine the relationships between visible reflectance, infrared emittance, and cloud-top temperatures for cirrus clouds.

#### 2. Data

Lidar backscatter data taken over Ft. McCoy (FMC), Wausau, and Madison, Wisconsin during the FIRE Cirrus IFO (October 19 - November 2, 1986) were used to define the cloud base, cloud-top altitude  $z_t$ , cloud thickness h, and effective cloud center altitude  $z_c$ . The last parameter is based on the distribution of backscatter intensity. Sassen et al. (1989) discuss the instrumentation and measurements in more detail. Soundings from Green Bay, Wisconsin, were used to determine the temperature-height relationships for all of the data. Cloud-top temperature  $T_t$  corresponds to  $z_t$  on the soundings. Mean cloud-center temperature  $T_c$  is found from  $z_c$ . temperatures were used to supplement the clear-sky temperatures  $T_{cs}$ derived from the satellite data. Half-hourly VIS  $(0.65\mu\text{m})$  and IR  $(11.5\mu\text{m})$ , 4-km data from the Geostationary Operational Environmental Satellite (GOES) were extracted along the wind vector at  $z_c$  for areas 4-pixels wide and ~ 16-pixels long. The VIS data were converted to reflectance  $\rho$ calibration of C. H. Whitlock (1988, personal communication). Clear-sky temperature is derived for each data strip whenever some clear pixels are observed in the VIS data. Otherwise, alternate methods are used to compute

 $T_{cs}$ . For all pixels (denoted as cloudy pixels) having an equivalent blackbody temperature, T <  $T_{cs}$  - 3K, the effective cloud beam emittance is

$$\epsilon = \{ [B(T) - B(T_{cs})] / [B(T_{z}) - B(T_{cs})] \},$$
 (1)

where  $T_z$  refers to the temperature at altitude z. The emittance is calculated twice for each satellite data strip using  $T_c$  and  $T_t$  in place of  $T_z$ . It is also assumed that

$$\epsilon = 1 - \exp(-\tau_e / \mu), \tag{2}$$

where the cloud IR optical depth is  $\tau_{\rm e}$  and  $\mu$  is cosine of the viewing zenith angle. The cloud reflectance,  $\rho_{\rm c}$ , is assumed to be related to the observed reflectance in the following manner.

$$\rho = T_{a} \rho_{c} + \rho_{s} T_{c} T_{u} + \alpha_{sd} (1 - \alpha_{d}) (1 - T_{c} - \alpha_{c}), \tag{3}$$

where  $T_a$  is the transmittance of the atmosphere above the cloud;  $T_c$  and  $T_u$  are the transmittances of the cloud to downward and upward direct radiation, respectively;  $\rho_s$  and  $\alpha_{sd}$  are the clear-sky bidirectional reflectance and diffuse albedos, respectively; and  $\alpha_d$  and  $\alpha_c$  are the albedos of the cloud to diffuse and total radiation, respectively. Thus,  $\alpha_d = \alpha_d(\tau_v)$ ;  $\alpha_c = \alpha_c(\tau_v, \mu_0)$ ; and  $\rho_c = \chi_c\alpha_c$ . The cosine of the solar zenith angle is  $\mu_0$ . It is assumed that the anisotropic reflectance correction factor  $\chi_c$  is independent of the VIS optical depth  $\tau_v$ . Values for  $\chi_c$  are taken from the empirical model of Minnis and Harrison (1984). The relationship bewtween  $\tau_v$  and cloud albedo is based on the results of Takano and Liou (1989) for a cloud containing randomly oriented hexagonal ice columns 125  $\mu$ m long and 50  $\mu$ m wide. Cloud transmittance also depends on the VIS optical depth. Mean values of T were computed for each discrete reflectance value.

In the ISCCP methodology, the scattering ratio, r, is used to relate the VIS optical depth to  $\tau_{\rm e}$  in order to correct for the semitransparency of the cloud. This quantity is equivalent to  $\tau_{\rm v}/\tau_{\rm e}$ . It is calculated here as

$$r = \sum_{i=1}^{N} (\tau_{v} / \tau_{e}) / N,$$

where N is the number of cloudy pixels for each data strip. Details concerning the data and methodology are given by Minnis et al. (1989).

### 3. Results and Discussion

Figure 1 shows an example of the observed values of  $\epsilon$ ,  $\rho$ , and  $\alpha$ cirrus clouds over FMC at 15 UTC, October 28, 1986. Averaging the temperatures minimizes some of the noise introduced by cloud structure variations. Mean values of r computed at each half hour for the cirrus days during the IFO (October 22, 27, 28, and 30; November 1 and 2) are summarized in Table 1. It is expected that r is constant for a given cirrus cloud. Thus, the systematic variation of r with scattering angle heta indicates that the values of  $\chi_c$  used here are inadequate. The mean r = 2.1, however, is very close to the value of 2.0 expected for large-particle clouds (e.g., Hansen and Travis, 1974). Assuming that the computed mean value of r is correct, it is possible to compute more realistic values of  $\chi_{\rm c}$  by forcing the results to yield r = 2.1 for each data strip. The average values of  $\chi_c$  resulting from this process are shown as the dashed lines in Fig. 2. Obviously, the cirrus clouds scatter radiation more anisotropically than the average cloud. Optically thin clouds tend to mimic the single-scattering phase function. Optically thick clouds tend to reduce anisotropy through multiple scattering. Scattering ratios from Table 1 averaged according to  $T_{\rm c}$  are shown in Fig. 3. The average values of r are 2.5  $\pm$  0.2 for T<sub>c</sub> < 235 K and 1.9  $\pm$  0.2 for T<sub>c</sub> > 235 K. The larger values of  $\, r \,$  are significantly different from 2.0 suggesting that the colder clouds contain a substantial number of small particles. The temperature dependency of r is similar to that found by Platt and Dilley (1981). The variation of cloud emittance (corrected for  $\mu$ ) with T given in Fig. 4 for the case study (October 27-28) is much like that determined by Platt et al. (1987) from ground observations. The data for the entire period, however, show greater increases of  $\epsilon$  with temperature than found by Platt et al. (1987). Cloud thickness (Fig. 5) during the case study is also similar to the Platt et al. (1987) results. Ratios of  $\epsilon(T_t) / \epsilon(T_c)$  are plotted against  $T_c$  in Fig. 6. The nearly linear dependency of these ratios on  $T_{\mathbf{c}}$  indicates that the actual cloudtop temperature can be determined if  $T_c$  is found first. Other results from these analyses (Minnis et al., 1989) indicate that cloud shadows have a significant effect on the retrieved cloud albedos. This shadowing can result in "dark" pixels which are colder than  $T_{\rm cs}$  and darker than  $\rho_{\rm s}$ . Their effects must be considered in the retrieval of cirrus properties from VIS and IR data.

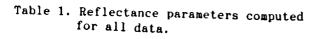
## 4. Concluding Remarks

The analysis of IFO data shows consistency with earlier cirrus studies suggesting that some simple parameterizations of cirrus cloud properties can be effectively utilized. Other results indicate that additional parameters may be derived from the satellite data. More cases taken over other areas will be needed to confirm the results presented here.

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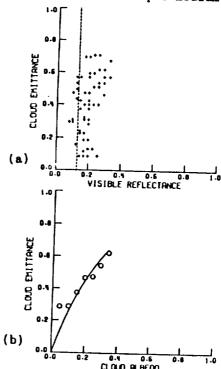


Fig. 1. Cloud emittances and observed (a) reflectances and (b) cloud albedos for cirrus clouds over FMC at 1500 UTC, October 28, 1986.

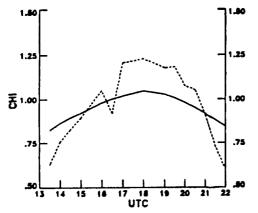


Fig. 2. Mean anisotropic reflectance correction factors for the IFO. The solid line is the nominal model; the dashed line represents the values derived for r = 2.1.

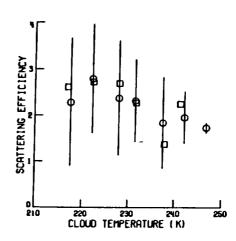


Fig. 3. Variation of mean scattering ratio with cloud-center temperature.

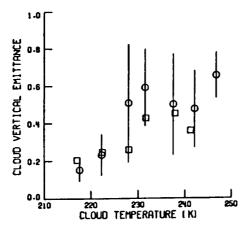


Fig. 4. Variation of cloud vertical beam emittance with cloud-center temperature.

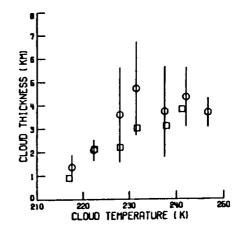


Fig. 5. Dependence of cloud thickness on cloud-center temperature.

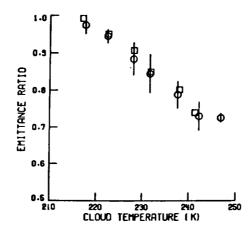


Fig. 6. Variation of emittance ratio with cloud-center temperature.